

$$Pv = R \cdot T$$

$$ds = \delta q + \delta L$$

$$a = \sqrt{\left(\frac{dp}{d\rho}\right)_{\text{scos}}} = \sqrt{\gamma RT}$$

$$V_{ac} = \sqrt{\frac{2 \cdot \varepsilon_t}{m}} = \sqrt{\frac{3 \cdot K_t}{m}} = \sqrt{\frac{3 \cdot R \cdot T}{M}}$$

$$\varepsilon_{int} = \frac{L}{2} \cdot RT$$

$$\bar{\varepsilon}_{\text{mediamol}} = \frac{L}{2} k_b T$$

$$\text{con } k_b = 1.38062 \cdot 10^{-23}$$

$$\beta = \frac{1}{V} \cdot \frac{dV}{dp}$$

Gradi di libertà			
	Traslazionali	Rotazionali	Vibrazionali
Monoatomiche	3	0	0
Biatomiche	3	2	2(3n-5)
Poliatomiche	3	3	2(3n-6)

$$C_p = \frac{L+2}{2} R \quad C_v = \frac{L}{2} R$$

$$x_i = \frac{N_i}{N_t} = \frac{P_i}{P_t} = \frac{V_i}{V_t} = \frac{\frac{\varepsilon_i}{m_i}}{\sum \left(\frac{\varepsilon_i}{m_i}\right)}$$

$$c_i = \frac{m_i}{m_t} = \frac{x_i N_i}{\sum x_i N_i}$$

	Unità mole	Unità massa
Energia	$\epsilon_{\text{misc}} = \sum x_i \epsilon_i$	$e_{\text{misc}} = \sum c_i \epsilon_i$
Entalpia	$h_{\text{misc}} = \sum x_i h_i$	$h_{\text{misc}} = \sum c_i h_i$
Cv	$C_{v\text{misc}} = \sum x_i C_{v_i}$	$C_{v\text{misc}} = \sum c_i C_v$
Cp	$C_{p\text{misc}} = \sum x_i C_{p_i}$	$C_{p\text{misc}} = \sum c_i C_p$
γ	$\gamma_{\text{misc}} = \frac{\sum x_i (L_i + 2)}{\sum x_i L_i}$	$\gamma_{\text{misc}} = \frac{C_{p\text{misc}}}{C_{v\text{misc}}}$

$$C_{p\text{mis}} = \sum C_{p_i} N_i$$

Aria standard	
T	15°=288 K
P	760mm HG =101325 Pa
ρ	1.225 kg/m ³

μ	$1.781 \cdot 10^{-5} \text{ kg/ms} = \text{Pas}$
ν	$1.454 \cdot 10^{-5} \text{ m}^2/\text{s}$

$$\mu = S \frac{T^{\frac{3}{2}}}{(\chi + T)} \quad \text{con } S = 1.46 \cdot 10^{-6} \text{ kg}/(\text{m s K}^{0.5}) \quad \text{e } \chi = 110\text{K} \quad \text{per aria}$$

$$\mu = \mu_{rif} * \left(\frac{T}{T_{rif}}\right)^w \quad \text{con } w = 0.75 \quad \text{per l'aria}$$

$$\mu_{mis} = \sum_{i=1}^n \mu_i \left(1 + \sum_{k=1, k \neq i}^n G_{ik} \frac{X_k}{X_i}\right)^{-1}$$

$$\mu_{mis} = \frac{\mu_1}{1 + G_{12} * \frac{X_2}{X_1}} + \frac{\mu_2}{1 + G_{21} * \frac{X_1}{X_2}} \quad \text{con} \quad G_{12} = \frac{\left(1 + \left(\frac{\mu_1}{\mu_2}\right)^{\frac{1}{2}} * \left(\frac{M_2}{M_1}\right)^{\frac{1}{4}}\right)^2}{2^{\frac{2}{3}} * \left(1 + \left(\frac{M_1}{M_2}\right)\right)^{\frac{1}{2}}}$$

$$ds = C_v * \frac{dT}{T} + \frac{R}{M} \frac{dV}{V}$$

$$C_p = \frac{\gamma}{\gamma - 1} * \frac{R}{M}$$

$$\frac{P_1}{P_2} = \left(\frac{T_1}{T_2}\right)^{\frac{\gamma}{\gamma - 1}}$$

$$\frac{\rho_1}{\rho_2} = \left(\frac{T_1}{T_2}\right)^{\frac{1}{\gamma - 1}}$$

Equazioni di continuit 

$$\frac{\partial}{\partial t} \int_{vol} \rho \, dvol = - \int_S \rho (\vec{V} \cdot \vec{n}) \, dS \quad \text{massa}$$

$$\frac{\partial}{\partial t} \int_{vol} \rho \vec{V} \, dvol = - \int_S \rho \vec{V} (\vec{V} \cdot \vec{n}) \, dS - \int_S p(\vec{n}) \, dS - \int_{vol} \rho \vec{F} \, dvol \quad \text{qdm}$$

$$\frac{\partial}{\partial t} \int_{vol} \rho \left(\epsilon + \frac{\vec{V}^2}{2}\right) \, dvol = - \int_S \rho \left(\epsilon + \frac{\vec{V}^2}{2}\right) (\vec{V} \cdot \vec{n}) \, dS + \int_{vol} \rho \dot{q} \, dvol + \dot{Q}_{visc} - \int_S p(\vec{V} \cdot \vec{n}) \, dS - \int_{vol} \rho (\vec{F} \cdot \vec{V}) \, dvol$$

energia

Forma conservativa

$$\frac{\partial p}{\partial t} + \frac{\partial \rho u}{\partial x} + \frac{\partial \rho v}{\partial y} + \frac{\partial \rho w}{\partial z} = 0 \quad \text{massa}$$

$$\frac{\partial \rho u}{\partial t} + \frac{\partial \rho u u}{\partial x} + \frac{\partial \rho v u}{\partial y} + \frac{\partial \rho w u}{\partial z} = - \frac{\partial p}{\partial x} + \rho f_x \quad \text{quantit  di moto lungo x}$$

$$\rho \frac{D\left(\epsilon + \frac{\vec{V}^2}{2}\right)}{Dt} = -\nabla(\rho \vec{V}) + \rho \dot{q} + \rho \vec{F} \cdot \vec{V} \quad \text{energia}$$

Equazioni di continuit 

Forma non conservativa

$$\frac{\partial \rho}{\partial t} + u \frac{\partial \rho}{\partial x} + v \frac{\partial \rho}{\partial y} + w \frac{\partial \rho}{\partial z} = 0 \quad \text{massa}$$

$$\frac{\partial \rho u}{\partial t} + u \frac{\partial \rho u}{\partial x} + v \frac{\partial \rho u}{\partial y} + w \frac{\partial \rho u}{\partial z} = -\frac{\partial p}{\partial x} + \rho f_x \quad \text{quantità di moto lungo x}$$

Grandezze totali

$$\rho_0 = \rho \left(\frac{T_0}{T} \right)^{\frac{1}{\gamma-1}} = \left(1 + \frac{\gamma-1}{2} M^2 \right)^{\frac{1}{\gamma-1}}$$

$$p_0 = p \left(\frac{T_0}{T} \right)^{\frac{\gamma}{\gamma-1}} = \left(1 + \frac{\gamma-1}{2} M^2 \right)^{\frac{\gamma}{\gamma-1}}$$

$$\frac{T_0}{T} = 1 + \frac{\gamma-1}{2} M^2$$

$$\frac{V_i^2}{2} = \frac{a_0}{\gamma-1}$$

Grandezze critiche

$$a_* = V_1 V_2$$

$$M_{1*} * M_{2*} = 1$$

$$M_{1*}^2 = \frac{1 + \frac{\gamma-1}{2} M_1^2}{\gamma M_1^2 - \frac{\gamma-1}{2}}$$

$$\frac{a_*^2}{a_0^2} = \frac{2}{\gamma+1} = \frac{T_*}{T_0}$$

$$\frac{P_*}{p_0} = \left(\frac{2}{\gamma+1} \right)^{\frac{\gamma}{\gamma+1}}$$

$$\frac{\rho_*}{\rho_0} = \left(\frac{2}{\gamma+1} \right)^{\frac{1}{\gamma-1}}$$

$$T_0 = \frac{\gamma+1}{\gamma-1} \frac{V_{cr}^2}{2c_p}$$

$$T \nabla S = \nabla H - \nabla \times \nabla \times \vec{V} \quad \text{teorema di crocco}$$

Relazioni di Rankine-Hugoniot

$$M_2^2 = \frac{1 + \frac{\gamma-1}{2} * M_1^2}{\gamma M_1^2 - \frac{\gamma-1}{2}}$$

$$\frac{P_2}{P_1} = 1 + \frac{2\gamma}{\gamma+1} (M_1^2 - 1)$$

$$\frac{\rho_2}{\rho_1} = \frac{v_1}{v_2} = \frac{(\gamma + 1)M_1^2}{2 + (\gamma - 1)M_1^2}$$

$$\frac{T_2}{T_1} = \left(1 + \frac{2\gamma}{\gamma + 1}((M_1^2 - 1))\right) \left(\frac{2 + (\gamma - 1)M_1^2}{(\gamma + 1)M_1^2}\right)^\gamma$$

$$\frac{p_{02}}{p_{01}} = \frac{\rho_{02}}{\rho_{01}} = \left(1 + \frac{2\gamma}{\gamma + 1}((M_1^2 - 1))\right)^{\frac{-1}{\gamma - 1}} * \left(\frac{2 + (\gamma - 1)M_1^2}{(\gamma + 1)M_1^2}\right)^{\frac{-\gamma}{\gamma - 1}}$$

$$S_2 - S_1 = C_v \ln \left(\left(1 + \frac{2\gamma}{\gamma + 1}((M_1^2 - 1))\right) \left(\frac{2 + (\gamma - 1)M_1^2}{(\gamma + 1)M_1^2}\right)^\gamma \right)$$

$$\frac{S_2 - S_1}{R^*} = -\ln \frac{p_{02}}{p_{01}} = \frac{\gamma}{\gamma - 1} \ln \frac{T_2}{T_1} - \ln \frac{P_2}{P_1}$$

Flusso di fanno

$$\frac{4FL^*}{D_i} = \left(\frac{1 - M^2}{\gamma M^2}\right) + \left(\frac{\gamma + 1}{2\gamma}\right) \ln \left(\frac{M^2}{\left(\frac{2}{\gamma + 1}\right) \left(1 + \frac{\gamma - 1}{2} M^2\right)} \right)$$

$$\frac{p}{p^*} = \frac{1}{M \sqrt{\left(\frac{2}{\gamma + 1}\right) \left(1 + \frac{\gamma - 1}{2} M^2\right)}}$$

$$\frac{\rho}{\rho^*} = \frac{1}{M} \sqrt{\left(\frac{2}{\gamma + 1}\right) \left(1 + \frac{\gamma - 1}{2} M^2\right)}$$

$$\frac{T}{T^*} = \frac{1}{\left(\frac{2}{\gamma + 1}\right) \left(1 + \frac{\gamma - 1}{2} M^2\right)}$$

$$\frac{V}{V^*} = \frac{M}{\sqrt{\left(\frac{2}{\gamma + 1}\right) \left(1 + \frac{\gamma - 1}{2} M^2\right)}}$$

$$\frac{P_2}{P_0} = \frac{1}{M} \left(\left(\frac{2}{\gamma + 1}\right) \left(1 + \frac{\gamma - 1}{2} M^2\right) \right)^{\frac{\gamma + 1}{2(\gamma - 1)}}$$

flusso di rayleigh

$$H_{01} + q = H_{02}$$

$$\frac{P_2}{P_1} = \frac{1 + \gamma M_1^2}{1 + \gamma M_2^2}$$

$$\frac{T_2}{T_1} = \frac{1 + M_1^2}{1 + M_2^2} \left(\frac{M_2}{M_1}\right)^2$$

$$\frac{\rho_2}{\rho_1} = \frac{1 + \gamma M_1^2}{1 + \gamma M_2^2} \left(\frac{M_2}{M_1}\right)^2$$

$$\frac{p_{02}}{p_{01}} = \frac{P_2}{P_1} * \left(\frac{1 + \frac{\gamma-1}{2} M_2^2}{1 + \frac{\gamma-1}{2} M_1^2} \right)^{\frac{\gamma}{\gamma-1}}$$

$$\frac{T_{02}}{T_{01}} = \frac{T_2}{T_1} * \left(\frac{1 + \frac{\gamma-1}{2} M_2^2}{1 + \frac{\gamma-1}{2} M_1^2} \right)$$

$$S_2 - S_1 = -\ln \frac{p_{02}}{p_{01}} = C_p \ln \left(\frac{1 + M_1^2}{1 + M_2^2} \left(\frac{M_2}{M_1} \right)^2 \right) - \frac{R}{M} \ln \left(\frac{1 + \gamma M_1^2}{1 + \gamma M_2^2} \right)$$

urtoobliquo

$$M_2^2 = \frac{1}{\sin(\beta - \theta)} * \frac{2 + (\gamma - 1) M_1^2 (\sin \beta)^2}{2\gamma M_1^2 (\sin \beta)^2 - (\gamma - 1)}$$

$$\tan \theta = \frac{2}{\tan \beta} * \left(\frac{M_1^2 (\sin \beta)^2 - 1}{M_1^2 (\gamma + \cos(2\beta)) + 2} \right)$$

$$v = \sqrt{\frac{\gamma+1}{\gamma-1}} \tan^{-1} \left(\sqrt{\frac{\gamma-1}{\gamma+1} * (M^2 - 1)} \right) - \tan^{-1} \sqrt{M^2 - 1}$$

Condotto senza riflessioni

$$\frac{l}{h} = M_1 \left(\frac{1 + \frac{\gamma-1}{2} M^2}{1 + \frac{\gamma-1}{2} M_1^2} \right)^{\left(\frac{\gamma}{\gamma-1} + \frac{1}{2} \right)}$$

Teoria PP

Equazione del potenziale di perturbazione linearizzata

$$(1 - M_\infty^2) \frac{\partial \phi^2}{\partial x} + \frac{\partial \phi^2}{\partial y} + \frac{\partial \phi^2}{\partial z} = 0$$

Coefficienti regime subsonico compressibile

$$C_{P_{inc}} = \frac{2u'}{V_\infty}$$

$$C_{P_{comp}} = C_{P_{inc}} \frac{1}{\beta} = C_{P_{inc}} \frac{1}{\sqrt{1 - M_\infty^2}}$$

$$C_{P_{comp}} = \frac{C_{P_{inc}}}{\sqrt{1 - M_\infty^2} + \left(\frac{M_\infty^2 \left(1 + \frac{\gamma+1}{2} M_\infty^2 \right)}{2\sqrt{1 - M_\infty^2}} \right) C_{P_{inc}}}$$

approssimazione di Laitone

$$C_{P_{comp}} = \frac{C_{P_{inc}}}{\sqrt{1-M^2_\infty} + \left(\frac{M^2_\infty}{1 + \sqrt{1-M^2_\infty}} \right) C_{P_{inc}}}$$

approssimazione di Karman Tsen

$$C_L = \int_0^l (C_{P_{ventre}} - C_{P_{dorso}}) \frac{dx}{c}$$

$$C_{L_{comp}} = C_{L_{inc}} \frac{1}{\beta} = C_{L_{inc}} \frac{1}{\sqrt{1-M^2_\infty}}$$

$$C_{M_{comp}} = C_{M_{inc}} \frac{1}{\beta} = C_{L_{inc}} \frac{1}{\sqrt{1-M^2_\infty}}$$

$$C_{P_{cr}} = \frac{2}{\gamma M^2_\infty} \left[\frac{1 + \frac{\gamma-1}{2} M^2_\infty}{1 + \frac{\gamma-1}{2}} \right]^{\frac{\gamma}{\gamma-1}}$$

$$C_{P_{inc_{min}}} = \sqrt{1-M^2_{cr}} \frac{2}{\gamma M^2_{cr}} \left[\left(\frac{1 + \frac{\gamma-1}{2} M^2_{cr}}{1 + \frac{\gamma-1}{2}} \right)^{\frac{\gamma}{\gamma-1}} - 1 \right]$$

Coefficienti regime supersonico compressibile linearizzato

$$C_P = \frac{2\theta}{\gamma} = \frac{2\theta}{\sqrt{M^2_\infty - 1}} \text{ con } \begin{cases} \theta < 0 \text{ ventre} \\ \theta > 0 \text{ dorso} \end{cases}$$

$$C_L = \frac{4(\alpha - \alpha_0)}{\sqrt{M^2_\infty - 1}}$$

$$C_D = \frac{4(\alpha - \alpha_0)^2}{\sqrt{M^2_\infty - 1}} + \frac{2}{c\sqrt{M^2_\infty - 1}} \int_0^c \left(\left(\frac{dy}{dx} \right)_i^2 + \left(\frac{dy}{dx} \right)_u^2 \right) dx$$

$$C_{M_\alpha} = -\frac{2(\alpha - \alpha_0)}{\sqrt{M^2_\infty - 1}} - \frac{2}{\sqrt{M^2_\infty - 1}} \frac{(S_u - S_l)}{c^2}$$

$$C_{M_0} = -\frac{2}{\sqrt{M^2_\infty - 1}} \frac{(S_u - S_l)}{c^2}$$

$$\frac{C_{L_{aia}}}{C_{L_{teorico}}} = 1 - \frac{1}{2A\sqrt{M^2_\infty - 1}}$$

$$Re = \frac{L \cdot V \cdot \rho}{\mu} = \frac{L \cdot V}{\nu}$$

$$Pr = \frac{\mu \cdot C_p}{\lambda} = \frac{2L + 4}{2L + 9} = \frac{4\gamma}{9\gamma - 5} \propto \frac{\delta_v}{\delta_\tau} \text{ per aria } = 0.71$$

Flusso di Couette

$$P_r \cdot \frac{u^2}{2} + h = const = P_r \cdot \frac{U_e^2}{2} + h_e$$

$$T_w = T_e + \frac{U_e^2}{2C_p} = T_\infty + \frac{U_\infty^2}{2C_p} = T_\infty \left(1 + \frac{\gamma-1}{2} M_\infty^2\right) \quad \text{Se Pr=1}$$

$$T = T_e + \frac{U_e^2}{2C_p} - \frac{u^2}{2C_p} = T_e + \frac{U_e^2}{2C_p} \left(1 - \frac{u^2}{U_e^2}\right) = T_e + \frac{U_e^2}{2C_p} (1 - \varphi^2) \quad \text{se adiabatico}$$

Strato limite termico

$$T_{rec} = T_e \left(1 + R \frac{\gamma-1}{2} M_e^2\right)$$

$$H_{rec} = h_e + R \frac{U_e^2}{2}$$

	Laminare	Turbolento	Couette
R	$\frac{1}{Pr^{1/2}}$	$\frac{1}{Pr^{1/3}}$	Pr
A	$\frac{2}{3}$	0.6	1

$$\frac{u_e}{U_e} = \frac{H - H_w}{H_e - H_w} = \frac{h + \frac{U^2}{2} - h_w}{h_e + \frac{U_e^2}{2} - h_w}$$

$$T = T_w(1 - \varphi) + T_\infty \varphi^2 + T_{rec}(\varphi - \varphi^2) \quad \text{profilo parabolico di T}$$

$$\frac{\tau_w(x)}{V_\infty} = - \frac{A \cdot q_w(x)}{h_\infty + R \frac{V_\infty^2}{2} - h_w} = \frac{A \cdot q_w(x)}{(T_{rec} - T_w)} \quad \text{Analogia di Reynolds}$$

$$K = \frac{c_D \cdot \rho_\infty \cdot V_\infty \cdot C_p}{2A} \quad \text{con } C_p \text{ calore specifico}$$

Metodo della sezione rappresentativa

$$T^* = T_w \cdot 0.5 + T_\infty \cdot 0.25 + T_{rec} \cdot 0.25$$

cd		
Laminare	$\frac{1.328}{\sqrt{R_{e=}}}$	$R_e < 5 \cdot 10^5$
Turbolento	$\frac{0.074}{R_{e=}^{1/4}}$	$R_e > 5 \cdot 10^5$
Alto turbolento	$\frac{0.0295}{R_{e=}^{1/7}}$	

$$C_{D_{incomp}} = \frac{\overline{\tau_w}_{incomp}}{\frac{1}{2} \rho_\infty V_\infty^2}$$

$$\overline{\tau_w}_{incomp} = C_{D_{incomp}} \frac{1}{2} \rho_\infty V_\infty^2$$

$$C_{D_{comp}} = \chi * C_{D_{incomp}}$$

$$\frac{\mu^*}{\mu_{\infty}} = \left(\frac{T^*}{T_{\infty}}\right) \quad \frac{\rho^*}{\rho_{\infty}} = \left(\frac{T^*}{T_{\infty}}\right)^{-1}$$

$$\chi = \left(\frac{\rho^* \mu^*}{\rho_{\infty} \mu_{\infty}}\right)^{\frac{1}{2}} \quad \chi = \left(\frac{T^*}{T_{\infty}}\right)^{\frac{w-1}{2}} \quad \text{per aria} \quad \chi = \left(\frac{T^*}{T_{\infty}}\right)^{-0.125} \quad \text{con} \quad \frac{T^*}{T_{\infty}} = 0.5 \frac{T_w}{T_{\infty}} + 0.25 + 0.25 \frac{T_{rec}}{T_{\infty}}$$

$$\chi_{Lam} = \left(\frac{T_w + T_{\infty}}{2T_{\infty}} + R \frac{\gamma - 1}{2} \frac{M_{\infty}}{4}\right)^{-0.125} \quad \text{per aria flusso laminare}$$

$$\chi_{Lam} = \left(1 + R \frac{\gamma - 1}{2} * \frac{3M_{\infty}}{4}\right)^{-0.125} \quad \text{se adiabatico}$$

$$\chi_{Turbo} = \left(\frac{T_w + T_{\infty}}{2T_{\infty}} + R \frac{\gamma - 1}{2} \frac{M_{\infty}}{4}\right)^{-0.65} \quad \text{per aria flusso turbolento}$$

$$\chi_{Turbo} = \left(1 + R \frac{\gamma - 1}{2} * \frac{3M_{\infty}}{4}\right)^{-0.65}$$

Cono

$$C_f(x)_{cono} = \sqrt{3} C_f(x)_{p.p.} \quad \text{Lamiare}$$

$$C_f(x)_{cono} = (1.1:1.5) C_f(x)_{p.p.} \quad \text{Turbolento}$$

Metodi integrali

$$\delta^* = \int_0^{\infty} \left(1 - \frac{u}{U_e}\right) dy \quad \text{Spessore di spostamento}$$

$$\theta^* = \int_0^{\infty} \frac{u}{U_e} \left(1 - \frac{u}{U_e}\right) dy \quad \text{Spessore di quantità di moto}$$

$$H = \frac{\delta^*}{\theta^*} = \begin{matrix} 2.5 \text{ laminare} \\ 1.3 \text{ turbolento} \end{matrix}$$

Metodo di Thwaites (laminare noto U_e)

$$\theta^2 = \frac{0.45\nu}{(U_e)^6} \int_0^x (U_e)^5 dx \quad \theta^2 = \frac{0.45\nu}{6k} \quad \text{nel punto d'arresto}$$

$$\lambda = \frac{\theta^2}{\nu} \left(\frac{dU_e}{dx}\right) \quad \lambda = 0.075 \quad \text{nel punto di arresto}$$

$$l = \left(\frac{du}{dy}\right)_{y=0} * \frac{\theta}{U_e(x)}$$

$$l = 0.22 + 1.57\lambda - 1.8\lambda^2 \quad \text{per flussi accelerati} \quad l = 0.328 \quad \text{nel punto di arresto}$$

$$l_{pa} = \frac{\tau_w}{\mu} \sqrt{\frac{0.45 \nu}{6k} \frac{1}{kx}} = 0.328$$

Cilindro

$$U_s \approx 2V_{\infty} \sin(\varphi) \quad \text{con} \quad \varphi = \frac{x}{R} \quad \text{ascissa curvilinea}$$

$$K = \left(\frac{dU_s}{dx} \right)_{x=0} = \frac{4V_{\infty}}{D}$$

$$\tau_{w, \text{incomp}} = 1.232 \sqrt{\rho_{\infty} \mu_{\infty} U_s^2 \left(\frac{dU_s}{dx} \right)}$$

$$q_w = 0.57 P_r^{-0.6} (\rho_s \mu_s K)^{\frac{1}{2}} \left(\frac{\rho_w \mu_w}{\rho_s \mu_s} \right)^{0.1} (H_s - h_w)$$

Sfera

$$U_s \approx \frac{3}{2} V_{\infty} \sin(\varphi) \quad \text{con} \quad \varphi = \frac{x}{R} \quad \text{ascissa curvilinea}$$

$$K = \left(\frac{dU_s}{dx} \right)_{x=0} = \frac{3V_{\infty}}{D}$$

$$\tau_{w, \text{incomp}} = 1.312 \sqrt{\rho_{\infty} \mu_{\infty} U_s^2 \left(\frac{dU_s}{dx} \right)}$$

$$q_w = 0.736 P_r^{-0.6} (\rho_s \mu_s K)^{\frac{1}{2}} \left(\frac{\rho_w \mu_w}{\rho_s \mu_s} \right)^{0.1} (H_s - h_w)$$

$$\frac{KD}{V_{\infty}} = \sqrt{8 \frac{\rho_{\infty}}{\rho_s}}$$

$$\frac{\rho_{\infty}}{\rho_s} = \frac{(\gamma + 1)M_{\infty}^2}{(\gamma + 1)M_{\infty}^2 + 2} \left[1 + \frac{\gamma + 1}{2} \frac{(\gamma - 1)M_{\infty}^2 + 2}{2\gamma M_{\infty}^2 - \gamma + 1} \right]^{\frac{1}{\gamma - 1}}$$

Metodo approssimato

$$\tau_{w, \text{comp}} = c \sqrt{\rho^* \mu^* U_s(x)^2 \left(\frac{dU_s}{dx} \right)}$$

$$q_w(x) = -\frac{c\sqrt{\rho^*\mu^*\left(\frac{dU_s}{dx}\right)}}{A} C_p(T_0 - T_w)$$

$$q_w(0) = -\frac{c\sqrt{\rho^*\mu^*K}}{P_r^{0.6}} C_p(T_0 - T_w)$$

Formula di Roming

$$q_w(0) = 0.0145 M_\infty^{3.1} \sqrt{\frac{P_\infty}{R}} \quad \text{risultato in } \frac{BTU}{ft^2s}$$

$$q_w(0) = 164.6724 M_\infty^{3.1} \sqrt{\frac{P_\infty}{R}}$$